Triaxial Deformation and Asynchronous Rotation of Rocky Planets in the Habitable Zone of Low-Mass Stars

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Rocky planets close to their host stars may be locked into spin-orbit resonances.



## Implications for Habitability



Two torques acting on planet's spin angular momentum near 3:2 spin-orbit resonance:

$$|\mathbf{T}_{\mathrm{tri}}| \simeq rac{21}{4} \epsilon le\left(rac{GM_{\star}}{a^3}
ight), \quad |\mathbf{T}_{\mathrm{tide}}| \simeq rac{3GM_{\star}^2R^5}{a^6}rac{k_2}{Q},$$

where  $\epsilon = (I_{yy} - I_{xx})/I_{zz}$ , with  $I_{zz} \ge I_{yy} \ge I_{xx}$ . For spin-orbit resonant capture, require  $|\mathbf{T}_{tril}| > |\mathbf{T}_{tide}|$ :

$$\Rightarrow \epsilon > 1.0 \times 10^{-6} \left(\frac{M_{\star}}{0.3 M_{\odot}}\right) \left(\frac{a}{0.1 \,\mathrm{AU}}\right)^{-3} \\ \times \left(\frac{\rho}{6 \,\mathrm{g/cm^3}}\right)^{-1} \left(\frac{k_2/Q}{10^{-3}}\right) \left(\frac{e}{0.01}\right)^{-1}.$$
(1)

Lower bound on  $\epsilon = (I_{yy} - I_{xx})/I_{zz}$  set by tides. What is upper bound on  $\epsilon$ , set by Geophysics? May super-earths achieve values of  $\epsilon$  capable of overcoming the lower limit set by tides? The two stresses on the planet's core are

$$|\mathbf{T}_{\text{grav}}| \sim \rho g R \epsilon, \quad |\mathbf{T}_{\text{elast}}| \sim \mu u.$$

 $\mu =$  shear modulus, u = strain. Requiring  $|\mathbf{T}_{elast}| \approx |\mathbf{T}_{grav}|$  gives

$$u \sim rac{
ho g R \epsilon}{\mu}.$$

Von-Mises criterion: Crust yields when  $u \ge u_{\rm crit} \sim 10^{-5}$ . In terms of ellipticity:

$$\epsilon \geq \epsilon_{\max} \sim \left(\frac{\mu}{\rho g R}\right) u_{\mathrm{crit}}.$$

Realistic model adds additional  $\sim 1/8$  geometric factor.

Assumptions:

- Constant density core  $\rho=\rho_{\rm c}$
- Homogeneous  $\mu = \text{constant}$
- Incompressible
- Fluid envelope with isothermal or constant density equation of state

## Bare Core



$$\bar{u} \sim \left(\frac{\mu}{\rho g R \epsilon}\right) u.$$

$$ar{u}_{
m peak} = \max(ar{u}) = 0.195,$$
  
 $\Rightarrow u_{
m peak} \simeq rac{1}{7.9} rac{
ho g R}{\mu} \epsilon.$ 

Setting  $u_{\text{peak}} = u_{\text{crit}}$ :

$$\begin{split} \epsilon_{\rm max} &= 1.9 \times 10^{-5} \left( \frac{\mu}{10^{12}\,{\rm dyn/cm^2}} \right) \\ &\times \left( \frac{\rho}{6\,{\rm g/cm^3}} \right)^{-2} \left( \frac{R}{R_\oplus} \right)^{-2} \left( \frac{u_{\rm crit}}{10^{-5}} \right). \end{split}$$

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 $\epsilon_{max}$  for constant density ocean.

Thin Isothermal atmosphere, equation of state  $\rho = p/c_{\rm s}^2$ ,  $c_{\rm s}^2 \ll g_{\rm c}R_{\rm c}$ .

 $\epsilon_{
m max} = \left(1-rac{
ho_{
m a}}{
ho_{
m c}}
ight)^{-1}\epsilon_{
m max,\ bare\ core}.$ 

Constant density ocean  $\rho = \rho_0$ :

$$\epsilon_{\max} = H\left(rac{R_{
m c}}{R},rac{
ho_{
m o}}{
ho_{
m c}}
ight)\epsilon_{
m max,\ bare\ core}.$$

Body	Observed $\epsilon$	$\epsilon_{\sf max}$
Mercury	$1.3 imes10^{-4}$	$1.6 imes10^{-4}$
Venus	$5.4 imes10^{-6}$	$2.9 imes10^{-5}$
Earth	$1.5 imes10^{-5}$	$2.3 imes10^{-5}$
Mars	$5.5 imes10^{-4}$	$1.6 imes10^{-4}$
Moon	$2.3 imes10^{-4}$	$8.7 imes10^{-4}$

Observed  $\epsilon$  computed through (Yoder 1995)

$$\epsilon \approx \frac{I_{yy} - I_{xx}}{2MR^2/5} = 10C_{22},$$

where  $C_{22}$  is gravitational coefficient.

• Maximal Triaxiality of rocky planet given by

$$\begin{split} \epsilon_{\rm max} &= 1.9 \times 10^{-5} \left( \frac{\mu}{10^{12}\,{\rm dyn/cm^2}} \right) \left( \frac{\rho}{6\,{\rm g/cm^3}} \right)^{-2} \\ & \times \left( \frac{R}{R_\oplus} \right)^{-2} \left( \frac{u_{\rm crit}}{10^{-5}} \right). \end{split}$$

- Fluid Envelope generally works to decrease  $\epsilon_{max}$ .
- Maximal triaxialities of Solar System Bodies close to observed values.